Rotating version of the static R=0 Lorentzian wormhole: geometry, matter and shadow



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A static wormhole with zero Ricci scalar (R=0)



$$ds^2=-(p+1)^{-2}\left(p+\sqrt{1-rac{2M}{r}}
ight)^2dt^2+rac{dr^2}{1-rac{2M}{r}}+r^2\left(d heta^2+\sin^2 heta d\phi^2
ight)$$

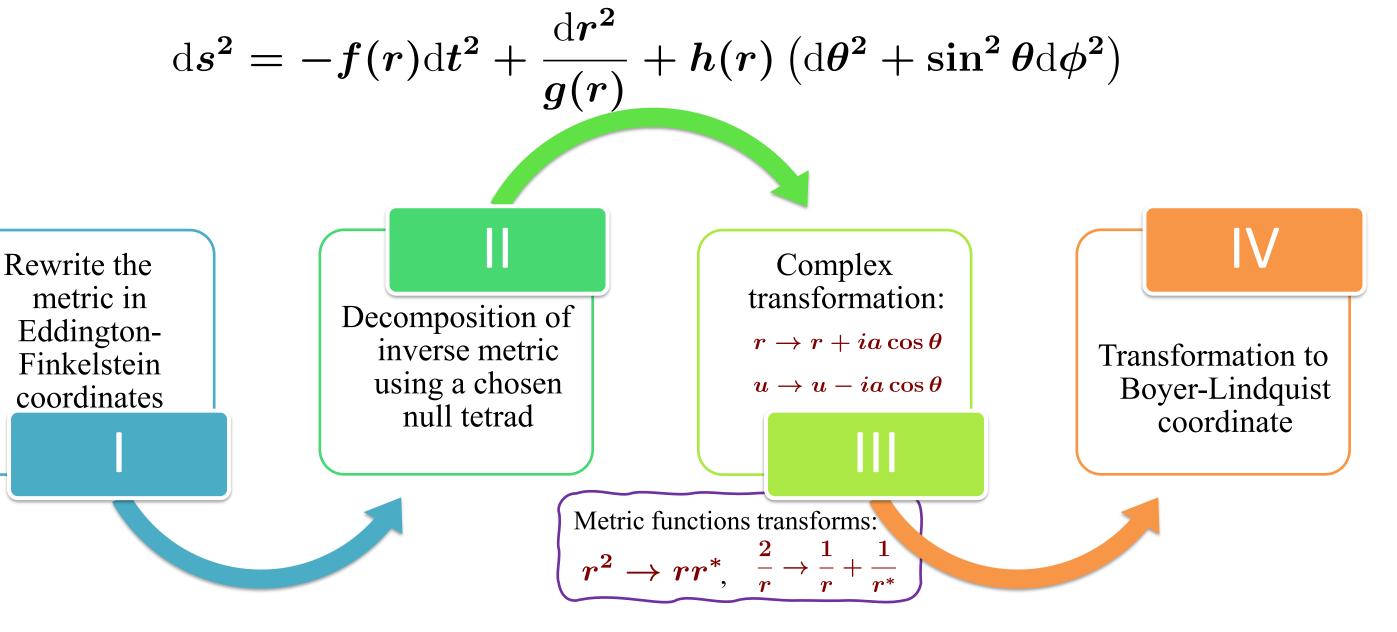
Ricci scalar, R=0

Static, spherically symmetric and asymptotically flat Lorentzian wormhole

Required matter violates energy conditions (assuming GR)

Phys. Rev. D 65, 064004 (2002), Phys. Rev. D 65, 084040 (2002).

Newman-Janis (NJ) algorithm and its failure



Step IV fails: No global coordinate transformation exists for the given metric.

J. Math. Phys. (N.Y.) 6, 915 (1965).

The Azreg-Ainou (AA) method

First two steps are same as NJ algorithm.

Difference comes in the step III

Assumptions under complex transformation: No need of complexification $f(r),g(r),h(r) o A(r, heta,a),B(r, heta,a),\Psi(r, heta,a)$ $\lim_{a o 0} A(r, heta,a) = f(r), \quad \lim_{a o 0} B(r, heta,a) = g(r), \quad \lim_{a o 0} \Psi(r, heta,a) = h(r).$

- \circ The step IV is performed through: $\mathrm{d} u = \mathrm{d} t \frac{h\sqrt{g/f + a^2}}{gh + a^2}\mathrm{d} r, \qquad \mathrm{d} \phi = \mathrm{d} arphi \frac{a}{gh + a^2}\mathrm{d} r$
- The three unknows (A, B, Ψ) of the new rotating metric fixed through:

$$g_{tr}=0, \qquad g_{r\phi}=0 \quad ext{and} \quad G_{r\theta}=0$$

 \circ At, a o 0, $ds_{rotating}^2 = \sqrt{g/f} ds_{static}^2$ (extra conformal factor) Phys. Letts. B 730, 95–98 (2014).

The new rotating wormhole using AA method

- o To remove the conformal factor, we start with the following transformed static metric: $ds^2 = -f(\ell)dt^2 + \frac{d\ell^2}{f(\ell)} + h(\ell)\left(d\theta^2 + \sin^2\theta d\phi^2\right)$
- \circ Applying all steps of AA method, we have the rotating metric in ℓ , converting ℓ to r,

$$\Delta(r) = \left(1-rac{2M}{r}
ight) \left[rac{\left(p+\sqrt{1-2M/r}
ight)^2r^2+a^2(p+1)^2}{\left(p+\sqrt{1-2M/r}
ight)^2}
ight]$$

 $p = 0 \rightarrow Kerr black hole$

 $a = 0 \rightarrow static wormhole$

$$m(r) = rac{r}{2} \left[1 - rac{\left(p + \sqrt{1 - 2M/r}
ight)^2}{(p+1)^2}
ight]
otag \ \Sigma = r^2 + a^2 \cos^2 heta$$

$$m(r) = \frac{r}{2} \left[1 - \frac{\left(p + \sqrt{1 - 2M/r} \right)^2}{(p+1)^2} \right] \\ \Sigma = r^2 + a^2 \cos^2 \theta$$

$$ds^2 = -\left[1 - \frac{2m(r)r}{\Sigma} \right] dt^2 - \frac{4m(r)ra\sin^2 \theta}{\Sigma} dt d\varphi + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 + \sin^2 \theta \left[r^2 + a^2 + \frac{2m(r)ra^2 \sin^2 \theta}{\Sigma} \right] d\varphi^2$$

The shape function satisfies all the requirements of wormhole

The required matter violates the energy conditions (assuming GR)

All independent curvature scalars are nonsingular everywhere.

Phys. Rev. D 111, 064010 (2025)

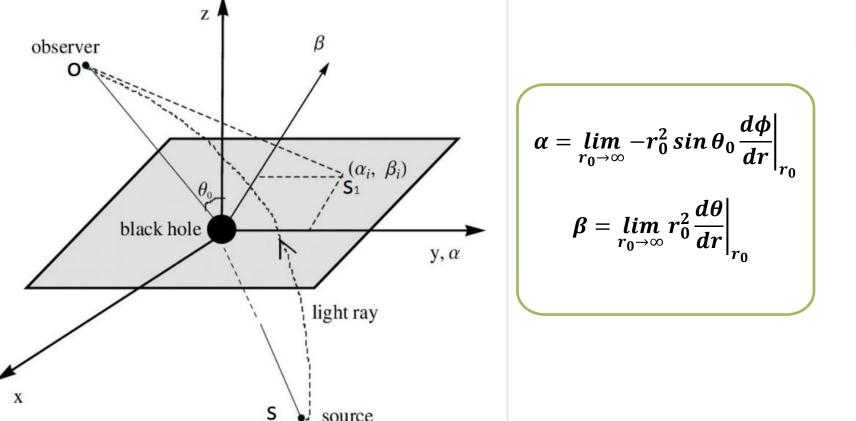
Shadow profiles and observations

Wormhole shadow **Unstable Spherical Orbits** or scatter. Vormhole's Throa **Entering Photon** Photon tragectory on the other side region of the Wormhole

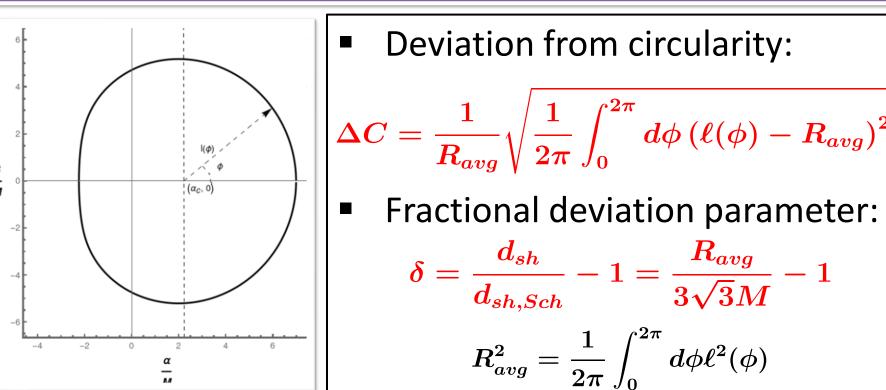
Celestial coordinates

- Photons can fall in the wormhole
- The boundary between infalling and outgoing is photon sphere.

Theoretical shadow profiles --p = 0.0-- a = 0.5 M - a = 1.0 M -p = 0.5-p = 1.0 β/M p = 0.5 and $\theta_0 = 90^{\circ}$ a = 0.99 M and θ_0 = 90°



Observable quantities

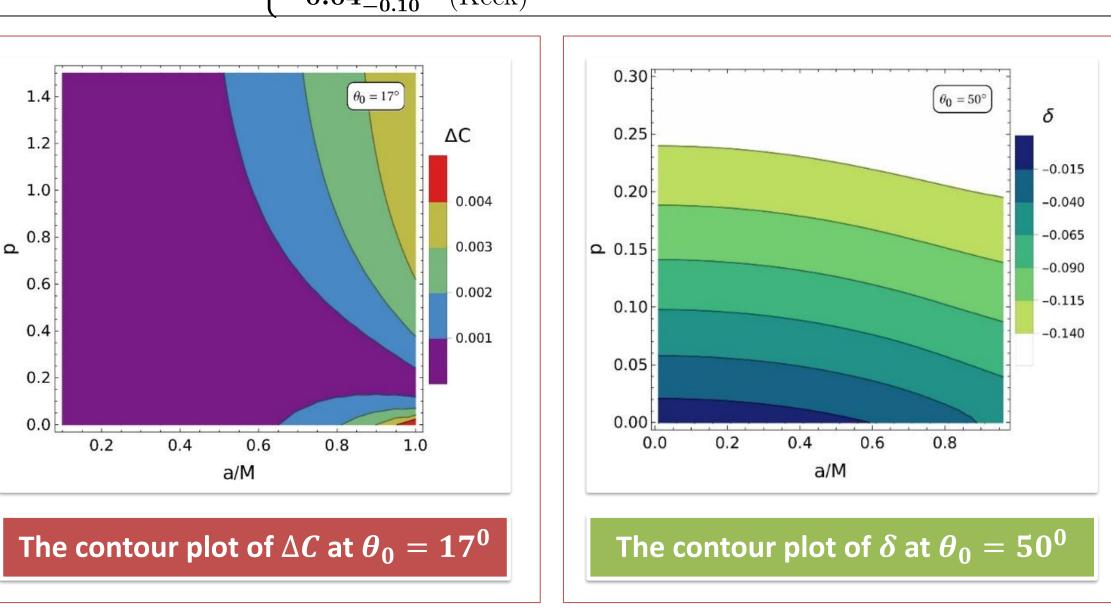


EHT results and comparison with the metric

- The observed angular diameter for M87* is $42 \pm 3 \mu as$, and for SgrA* it is $51.8 \pm 2.3 \mu as$. APJ. Lett. 875, L1 (2019), APJ. Lett. 930, L12 (2022)
- For M87* $\Delta C \lesssim 10\%$

APJ. Lett. 875, L5 (2019)

- - APJ. Lett. 930 L17 (2022)



- All values of p respect the restriction from M87*
- \circ p < 0.24 is required for Sgr A*